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INSTANTONS ON FLAT SPACE: EXPLICIT CONSTRUCTIONS

JASON D. LOTAY AND THOMAS BRUUN MADSEN

ABSTRACT. In this note, we revisit some well-known examples of instantons on flat space that were originally discovered in the physics literature. In particular, we explain how the basic instanton on \mathbb{R}^4 , with its flat hyperkähler structure, has natural generalisations to \mathbb{R}^7 and \mathbb{R}^8 viewed as flat G_2 - and $\text{Spin}(7)$ -manifolds, respectively. We also provide the details of an arguably less known construction of ASD instantons on \mathbb{H}^n , in the sense of quaternionic geometry.

1. INTRODUCTION

The idea of generalising the concept of (anti-)self-duality to higher dimensions has been around for about four decades, gaining increased attention with the advances in special and exceptional holonomy (see, e.g., [14] for a recent collection of surveys) and the potential link between higher-dimensional gauge theory and enumerative invariants, initiated in the seminal papers [6, 5]. Instantons in higher dimensions, which are connections whose curvature satisfy these generalised anti-self-duality conditions, have been explored both on non-compact and compact spaces, with most of the current interest centering around G_2 -instantons, e.g., [16, 17, 20, 21, 22].

The purpose of this note is to revisit some of the early constructions of instantons that appeared in the physics literature [1, 8, 13, 11], rephrasing explicit constructions in a way that highlights similarities and contextualising them in view of more recent developments. In particular, we shall explain that the basic (or BPST) instanton on \mathbb{R}^4 , with its flat hyperkähler structure, has natural generalisations to \mathbb{R}^7 and \mathbb{R}^8 viewed as flat G_2 - and $\text{Spin}(7)$ -manifolds, respectively. This generalization is achieved via an ansatz based on the groups associated with the ambient flat structure. Surprisingly, however, this natural ansatz that we present does not produce instantons for \mathbb{R}^n equipped with the remaining Ricci-flat geometries from Berger's list [2].

1.1. G-instantons. To set the scene, consider having a principal K -bundle $P \rightarrow M$ over an oriented Riemannian n -manifold M . Given a connection form $\omega \in \Omega^1(P; \mathfrak{k})$ the associated curvature will be $\Omega \in \Omega^2(M; \mathfrak{k}_P)$, where \mathfrak{k} is the Lie algebra of K . When M comes equipped with a G -structure, for some $G \subset \text{SO}(n)$, then since we can identify the 2-forms on M at each point with the Lie algebra of $\text{SO}(n)$ we can decompose the bundle of 2-forms as

$$\Lambda^2 T^* M \cong \mathfrak{so}(n) \cong \mathfrak{g} \oplus \mathfrak{g}^\perp,$$

where the fibres of \mathfrak{g} are given by the Lie algebra of G . This splitting gives us a way of distinguishing connections that are particularly adapted to the geometry (cf. [19]).

Definition 1.1. A connection ω on P is called a *G-instanton* if the 2-form part of its curvature Ω takes values in the subbundle $\mathfrak{g} \subset \Lambda^2 T^* M$.

A natural setting where we have a distinguished G -structure on M is when the metric on M has special holonomy G (and thus the G -structure is torsion-free). Of course, holonomy reduction is trivially obtained on flat space $M = \mathbb{R}^n$ which is our focus in this note. Taking the group G in question to be $\text{SU}(2)$, G_2 and $\text{Spin}(7)$ in dimensions 4, 7 and 8 leads to the notions of (anti-)self-dual, G_2 - and $\text{Spin}(7)$ -instantons, respectively.

1.2. Results. The main observation of this note is that there is an ansatz which gives a unified perspective on the construction of the basic G-instantons on flat spaces.

Let $P_G \cong G \times \mathbb{R}^n$ be the reduction of the principal frame bundle of \mathbb{R}^n corresponding to a choice of G-structure on \mathbb{R}^n . We construct a family of connections on P_G as follows.

Away from the origin, we have the usual identification $\mathbb{R}^n \setminus \{0\} \cong \mathbb{R}_+ \times \mathcal{S}^{n-1}$, with radial coordinate r of \mathbb{R}_+ and we write $\partial_r = \partial/\partial r$. Let e_j , for $j = 1, \dots, \dim G$, define a Killing orthonormal basis of the Lie algebra \mathfrak{g} of G inside $\mathfrak{so}(n)$ and denote by β_j the elements defining a corresponding basis for \mathfrak{g} inside $\Lambda^2(\mathbb{R}^n)^* \cong \mathfrak{so}(n)$. Then define 1-forms α_j on \mathcal{S}^{n-1} via

$$\alpha_j = \partial/\partial r \lrcorner \beta_j|_{\mathcal{S}^{n-1}}.$$

We let $\omega \in \Omega^1(P_G; \mathfrak{g})$ be of the form

$$\omega = A(r) = ra(r) \sum_{j=1}^{\dim G} \alpha_j \otimes e_j \quad (1.1)$$

for some function $a : \mathbb{R}_+ \rightarrow \mathbb{R}$, where we may view ω as a family of connections $A(r)$ on $G \times \mathcal{S}^{n-1}$. We have added the factor of r in (1.1) for convenience.

We then show the following.

Theorem 1.2. *For $G = \mathrm{SU}(2), G_2, \mathrm{Spin}(7)$ for $n = 4, 7, 8$ respectively, the ansatz (1.1) yields a G-instanton on \mathbb{R}^n with structure group G for*

$$a(r) = c_n \frac{r}{r^2 + C},$$

where $C > 0$ is a constant and c_n is a fixed constant only depending on n . These are the basic G-instantons on \mathbb{R}^n for $n = 4, 7, 8$.

As observed in Proposition 2.4, in the 4-dimensional case, these instanton solutions are the only globally defined Yang-Mills connections satisfying our rotationally symmetric ansatz (1.1).

Remark 1.3. Given the result above, it is natural to ask what happens for the other Ricci-flat geometries, associated with groups $G = \mathrm{SU}(n)$, $n > 2$, on $\mathbb{R}^{2n} = \mathbb{C}^n$ and $\mathrm{Sp}(n)$, $n > 1$, on $\mathbb{R}^{4n} = \mathbb{H}^n$. However, direct computations show that the ansatz (1.1) does not produce G-instantons in these situations. This is perhaps another manifestation of the “exceptional” nature of the holonomy groups G_2 and $\mathrm{Spin}(7)$, as well as the special case of 4 dimensions where the group is $\mathrm{SU}(2) = \mathrm{Sp}(1)$.

We also discuss another construction of instantons on flat space which is related to Nahm’s equations and quaternionic structures. For any solution of Nahm’s equations, we obtain families of $\mathrm{Sp}(n)$ -instantons on \mathbb{H}^n , as explained by Proposition 3.6.

Remark 1.4. Throughout the article we will use (x_1, \dots, x_n) for coordinates on \mathbb{R}^n and let $dx^{ij\dots k}$ denote $dx^i \wedge dx^j \wedge \dots \wedge dx^k$. We also let $E_{k\ell}$ be the elementary $n \times n$ matrix with 1 in the (k, ℓ) entry and zero for all other entries.

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2. BASIC INSTANTONS

We will now show explicitly how the ansatz (1.1) recovers the basic instanton on \mathbb{R}^4 (with the flat metric) and leads to generalisations on \mathbb{R}^7 (with a flat G_2 -structure) and \mathbb{R}^8 (with a flat $\mathrm{Spin}(7)$ -structure).

2.1. On \mathbb{R}^4 . We first consider the construction of the basic instanton on the trivial bundle $SU(2) \times \mathbb{R}^4$ and so need the standard $SU(2)$ -structure on $\mathbb{R}^4 \cong \mathbb{C}^2$. This structure can be specified in terms of a non-degenerate 2-form σ and a $(2,0)$ -form Ψ on \mathbb{R}^4 given by:

$$\sigma = dx^{12} + dx^{34}, \quad \Psi = \psi + i\hat{\psi} = (dx^1 + idx^2) \wedge (dx^3 + idx^4).$$

With a view to later constructions, we can see this in terms of structures on $\mathbb{R}^4 \setminus \{0\} \cong \mathbb{R}_+ \times \mathcal{S}^3$, we have

$$\sigma = dr \wedge r\eta + \frac{r^2}{2}d\eta, \quad \Psi = (dr + ir\alpha) \wedge r\Phi, \quad (2.1)$$

for a (contact) 1-form η and a $(1,0)$ -form Φ on \mathcal{S}^3 . For later use, it is convenient to introduce the notation $\Phi = \phi + i\hat{\phi}$ for the decomposition of the $(1,0)$ -form Φ into real and imaginary parts.

The condition for a 2-form β to be in $\mathfrak{su}(2) \subset \Lambda^2(\mathbb{R}^4)^*$ can be expressed as

$$\Psi \wedge \beta = 0 \quad \text{and} \quad \sigma \wedge \beta = 0. \quad (2.2)$$

Remark 2.1. It is elementary to check that the condition (2.2) on β is equivalent to demanding that β is anti-self-dual with respect to the standard orientation on \mathbb{R}^4 . The first equality in (2.2) says that β is of type $(1,1)$ for the standard complex structure on $\mathbb{R}^4 \cong \mathbb{C}^2$, and we note that, as β is real, it suffices to check that $\beta \wedge \psi = 0$. The second equality in (2.2) implies that β is “trace-free” and so orthogonal to σ inside $\Lambda^2(\mathbb{R}^4)^*$.

In order to proceed, let us fix a basis of $\mathfrak{su}(2) \leq \mathfrak{so}(4)$. Recalling the elementary matrices $E_{k\ell}$, an orthonormal basis, up to scale, of $\mathfrak{su}(2)$ is given by the elements:

$$\begin{aligned} e_1 &= \frac{1}{2}(-E_{12} + E_{21} + E_{34} - E_{43}), & e_2 &= \frac{1}{2}(-E_{13} + E_{31} + E_{42} - E_{24}), \\ e_3 &= \frac{1}{2}(-E_{14} + E_{41} + E_{23} - E_{32}). \end{aligned}$$

This basis satisfies the usual commutation relations $[e_i, e_j] = -2e_k$ for cyclic permutations (i, j, k) of $(1, 2, 3)$. Then $\mathfrak{su}(2) \subset \Lambda^2(\mathbb{R}^4)^*$ is spanned by the corresponding 2-forms

$$\beta_1 = dx^{12} - dx^{34}, \quad \beta_2 = dx^{13} - dx^{42}, \quad \beta_3 = dx^{14} - dx^{23}.$$

Now let ω be a connection given by (1.1) with curvature Ω . It follows by Definition 1.1 and (2.2) that the $SU(2)$ -instanton condition is given by

$$\Omega \wedge \sigma = 0 \quad \text{and} \quad \Omega \wedge \psi = 0.$$

Equivalently, viewing ω as family of connections $A(r)$ over \mathcal{S}^3 and recalling (2.1), we can express the instanton condition as

$$\begin{aligned} rA' \wedge d\eta &= -2F_A \wedge \eta, & F_A \wedge d\eta &= 0, \\ rA' \wedge \alpha \wedge \hat{\phi} &= F_A \wedge \phi, & F_A \wedge \hat{\phi} \wedge \alpha &= 0, \end{aligned} \quad (2.3)$$

where F_A is the curvature of A .

Given the symmetry of the problem, for computational simplicity we can consider the point $(1, 0, 0, 0)$ in \mathcal{S}^3 so that we may identify the unit radial vector field ∂_r with $\partial_1 = \partial/\partial x_1$. We then have that $\alpha_j = dx^{j+1}$ and

$$\beta_1|_{\mathcal{S}^3} = -dx^{34}, \quad \beta_2|_{\mathcal{S}^3} = -dx^{42}, \quad \beta_3|_{\mathcal{S}^3} = -dx^{23}.$$

Now at $(1, 0, 0, 0) \in \mathcal{S}^3$ the 2-form F_A is given by:

$$F_A = 2(ra(r)(1 + ra(r)) \sum_{j=1}^3 \beta_j|_{\mathcal{S}^3} \otimes e_j.$$

It is now easy to see that the instanton condition as expressed in (2.3) reduces to a single ODE:

$$\frac{da}{dr} = \frac{a}{r}(1 + 2ar). \quad (2.4)$$

Solving (2.4), we recover the basic instanton on \mathbb{R}^4 . In summary:

Proposition 2.2. *The basic instanton on \mathbb{R}^4 arises from the ansatz (1.1) with $G = \mathrm{SU}(2)$ and*

$$a(r) = -\frac{r}{r^2 + C},$$

where C is a positive constant. □

Remark 2.3. It follows from the above computations that $\int_{\mathbb{R}^4} \|\Omega\|^2 \mathrm{vol}_{\mathbb{R}^4}$ is finite. Note that the $\mathrm{SU}(2)$ -instanton condition is conformally invariant since it is the anti-self-duality condition by Remark 2.1. Since the L^2 -norm of Ω is also conformally invariant, it follows that stereographic projection will take the basic instanton on \mathbb{R}^4 and give an instanton on \mathcal{S}^4 with a removable point singularity.

It seems natural to also ask, more generally, what are the possible solution to the Yang–Mills equation, arising from our connection ansatz (1.1) in this setting; that is, connections which are critical points for the Yang–Mills functional. Elementary computations show that the Yang–Mills equations are given by the second order non-linear ODE:

$$r^2 \frac{d^2 a}{dr^2} + 3r \frac{da}{dr} - 4ra^2(3 + 2ra) - 3a = 0. \quad (2.5)$$

To better understand this second order ODE, it is useful make the change of variables:

$$u(r) = \frac{a(r)}{r}.$$

Note that, in terms of u , the instanton condition (2.4) reads:

$$\frac{du}{dr} = 2ru^2.$$

Motivated by this, let:

$$v(r) = r^5 \left(\frac{du}{dr} - 2ru^2 \right).$$

Using these variables, (u, v) , the Yang–Mills condition (2.5) becomes equivalent to the following system of first order ODEs:

$$\frac{du}{dr} = 2ru^2 + \frac{v}{r^5} \quad \text{and} \quad \frac{dv}{dr} = -4ruv. \quad (2.6)$$

Now, it is not difficult to see that the instanton solutions of Proposition 2.2 exhaust all globally defined Yang–Mills solutions.

Proposition 2.4. *The only globally defined Yang–Mills solutions of (2.5) are the instanton solutions of Proposition 2.2.*

Proof. The statement follows if we can show that a globally defined solution (u, v) of (2.6) necessarily has $v \equiv 0$. In this case, $a(r) = ru(r)$ solves the instanton equation (2.4). By uniqueness of solutions to (regular) initial value problems, $v \equiv 0$ holds if we can show that $v(r_0) = 0$ for some $r_0 > 0$.

To start with, we observe that regularity of the system (2.6), at zero, forces $v(0) = 0$. Next choose $\varepsilon > 0$ so that either $u(r) \geq 0$, or $u(r) \leq 0$, for all $[0, \varepsilon]$.

In the case where $u(r) \geq 0$, we see that for $r \in [0, \varepsilon]$, $v'(r)$ will either vanish or have the opposite sign of $v(r)$. Given that $v(0) = 0$, this is only possible if v vanishes identically. So, in this case, $v(r_0) = 0$ for some $r_0 > 0$, as required.

Similarly, if $u(r) \leq 0$ on $[0, \varepsilon]$, choose $\delta \in \{-1, 1\}$ so that $\delta v(r) \geq 0$ for small enough r . Let $C \geq 0$ be a constant so that $0 \leq -4ru \leq C$ on $[0, \varepsilon]$. Then, for small r , we have that:

$$\delta v(r) = \delta (v(r) - v(0)) = \delta \int_0^r \frac{dv}{ds} ds = \int_0^r (-4su) \delta v ds \leq C \left| \int_0^r \delta v ds \right|.$$

Hence, by Grönwall's inequality $\delta v \equiv 0$, so $v \equiv 0$, for small enough values of r . So again, we have that $v(r_0) = 0$ for some $r_0 > 0$.

In conclusion, we have shown that for any globally defined solution (u, v) of (2.6), $a(r) = ru(r)$ is a solution of the instanton condition (2.4), as required. \square

2.2. On \mathbb{R}^7 . We now turn to the construction of what we shall call the basic instanton on the trivial bundle $G_2 \times \mathbb{R}^7$. This requires a geometric description of the compact Lie group G_2 which may be given in terms of the 3-form φ on \mathbb{R}^7 given by

$$\varphi = dx^{123} + dx^1 \wedge (dx^{45} + dx^{67}) + dx^2 \wedge (dx^{46} + dx^{75}) - dx^3 \wedge (dx^{47} + dx^{56}), \quad (2.7)$$

as $\text{Stab}_{\text{GL}(7, \mathbb{R})}(\varphi) = G_2 \subset \text{SO}(7)$. Since φ determines the standard metric and orientation on \mathbb{R}^7 , we can also consider its dual 4-form:

$$*\varphi = dx^{4567} + dx^{23} \wedge (dx^{45} + dx^{67}) + dx^{31} \wedge (dx^{46} + dx^{75}) - dx^{12} \wedge (dx^{47} + dx^{56}). \quad (2.8)$$

Rephrasing in terms of the standard nearly Kähler structure $(\sigma, \Psi = \psi + i\hat{\psi})$ on the unit sphere $S^6 \subset \mathbb{R}^7$, where σ is a 2-form and Ψ is a $(3, 0)$ -form, we have

$$\varphi = dr \wedge r^2 \sigma + r^3 \psi, \quad *\varphi = r^3 \hat{\psi} \wedge dr + \frac{1}{2} r^4 \sigma^2. \quad (2.9)$$

on $\mathbb{R}_+ \times S^6$.

It is well-known that the condition for a 2-form β to be in $\mathfrak{g}_2 \subset \Lambda^2(\mathbb{R}^7)^*$ can be expressed as either of the equivalent conditions

$$*(\varphi \wedge \beta) = -\beta \quad \text{or} \quad *\varphi \wedge \beta = 0. \quad (2.10)$$

The first condition manifestly has echoes of the anti-self-duality condition in 4 dimensions.

For an G_2 -instanton of the form (1.1), as well as substituting $\beta = \Omega$ in (2.10), the instanton condition can also be rephrased as:

$$A' = *_r(F_A \wedge r^3 \psi), \quad (2.11)$$

subject to the initial constraint

$$F_A(r_0) \wedge \sigma(r_0)^2 = 0, \quad (2.12)$$

see, e.g., [15]. Here we are using the decomposition (2.9) and $*_r$ denotes the Hodge star on the 6-sphere of radius r .

As in the previous case, to make our construction explicit, we need to establish an identification of $\mathfrak{g}_2 \leq \mathfrak{so}(7)$. Using the elementary matrices E_{kl} , a basis of \mathfrak{g}_2 is given as follows:

$$\begin{aligned} e_1 &= \frac{1}{4}(E_{23} - E_{32} - E_{45} + E_{54}), & e_2 &= -\frac{1}{4\sqrt{3}}(E_{23} - E_{32} + E_{45} - E_{54} - 2E_{67} + 2E_{76}), \\ e_3 &= \frac{1}{4}(E_{13} - E_{31} + E_{46} - E_{64}), & e_4 &= -\frac{1}{4\sqrt{3}}(E_{13} - E_{31} - E_{46} + E_{64} - 2E_{57} + 2E_{75}), \\ e_5 &= -\frac{1}{4}(E_{12} - E_{21} + E_{47} - E_{74}), & e_6 &= \frac{1}{4\sqrt{3}}(E_{12} - E_{21} - E_{47} + E_{74} + 2E_{56} - 2E_{65}), \\ e_7 &= \frac{1}{4}(E_{15} - E_{51} - E_{26} + E_{62}), & e_8 &= \frac{1}{4\sqrt{3}}(E_{15} - E_{51} + E_{26} - E_{62} + 2E_{37} - 2E_{73}), \\ e_9 &= -\frac{1}{4}(E_{14} - E_{41} - E_{27} + E_{72}), & e_{10} &= -\frac{1}{4\sqrt{3}}(E_{14} - E_{41} + E_{27} - E_{72} - 2E_{36} + 2E_{63}), \\ e_{11} &= \frac{1}{4}(E_{17} - E_{71} + E_{24} - E_{42}), & e_{12} &= \frac{1}{4\sqrt{3}}(E_{17} - E_{71} - E_{24} + E_{42} - 2E_{35} + 2E_{53}), \\ e_{13} &= -\frac{1}{4}(E_{16} - E_{61} + E_{25} - E_{52}), & e_{14} &= -\frac{1}{4\sqrt{3}}(E_{16} - E_{61} - E_{25} + E_{52} + 2E_{34} - 2E_{43}). \end{aligned}$$

Correspondingly $\mathfrak{g}_2 \subset \Lambda^2(\mathbb{R}^7)^*$ has basis:

$$\begin{aligned} \beta_1 &= -\frac{1}{4}(dx^{23} - dx^{45}), & \beta_2 &= \frac{1}{4\sqrt{3}}(dx^{23} + dx^{45} - 2dx^{67}), \\ \beta_3 &= -\frac{1}{4}(dx^{13} + dx^{46}), & \beta_4 &= \frac{1}{4\sqrt{3}}(dx^{13} - dx^{46} - 2dx^{57}), \\ \beta_5 &= \frac{1}{4}(dx^{12} + dx^{47}), & \beta_6 &= -\frac{1}{4\sqrt{3}}(dx^{12} - dx^{47} + 2dx^{56}), \\ \beta_7 &= -\frac{1}{4}(dx^{15} - dx^{26}), & \beta_8 &= -\frac{1}{4\sqrt{3}}(dx^{15} + dx^{26} + 2dx^{37}), \\ \beta_9 &= \frac{1}{4}(dx^{14} - dx^{27}), & \beta_{10} &= \frac{1}{4\sqrt{3}}(dx^{14} + dx^{27} - 2dx^{36}), \\ \beta_{11} &= -\frac{1}{4}(dx^{17} + dx^{24}), & \beta_{12} &= -\frac{1}{4\sqrt{3}}(dx^{17} - dx^{24} - 2dx^{35}), \\ \beta_{13} &= \frac{1}{4}(dx^{16} + dx^{25}), & \beta_{14} &= \frac{1}{4\sqrt{3}}(dx^{16} - dx^{25} + 2dx^{34}). \end{aligned}$$

For simplification, we proceed by considering the point $x \in \mathcal{S}^6$ with $x_1 = 1$ and $x_j = 0$ for all other j . Then at x we have:

$$\begin{aligned} \alpha_1 &= 0, & \alpha_2 &= 0, & \alpha_3 &= -\frac{1}{4}dx^3, & \alpha_4 &= \frac{1}{4\sqrt{3}}dx^3, & \alpha_5 &= \frac{1}{4}dx^2, \\ \alpha_6 &= -\frac{1}{4\sqrt{3}}dx^2, & \alpha_7 &= -\frac{1}{4}dx^5, & \alpha_8 &= -\frac{1}{4\sqrt{3}}dx^5, & \alpha_9 &= \frac{1}{4}dx^4, & \alpha_{10} &= \frac{1}{4\sqrt{3}}dx^4, \\ \alpha_{11} &= -\frac{1}{4}dx^7, & \alpha_{12} &= -\frac{1}{4\sqrt{3}}dx^7, & \alpha_{13} &= \frac{1}{4}dx^6, & \alpha_{14} &= \frac{1}{4\sqrt{3}}dx^6. \end{aligned}$$

Similarly, at x , we have:

$$\begin{aligned} \beta_1|_{\mathcal{S}^6} &= -\frac{1}{4}(dx^{23} - dx^{45}), & \beta_2|_{\mathcal{S}^6} &= \frac{1}{4\sqrt{3}}(dx^{23} + dx^{45} - 2dx^{67}), \\ \beta_3|_{\mathcal{S}^6} &= -\frac{1}{4}dx^{46}, & \beta_4|_{\mathcal{S}^6} &= \frac{1}{4\sqrt{3}}(-dx^{46} - 2dx^{57}), \\ \beta_5|_{\mathcal{S}^6} &= \frac{1}{4}dx^{47}, & \beta_6|_{\mathcal{S}^6} &= \frac{1}{4\sqrt{3}}(dx^{47} - 2dx^{56}), \\ \beta_7|_{\mathcal{S}^6} &= \frac{1}{4}dx^{26}, & \beta_8|_{\mathcal{S}^6} &= -\frac{1}{4\sqrt{3}}(dx^{26} + 2dx^{37}), \\ \beta_9|_{\mathcal{S}^6} &= -\frac{1}{4}dx^{27}, & \beta_{10}|_{\mathcal{S}^6} &= \frac{1}{4\sqrt{3}}(dx^{27} - 2dx^{36}), \\ \beta_{11}|_{\mathcal{S}^6} &= -\frac{1}{4}dx^{24}, & \beta_{12}|_{\mathcal{S}^6} &= \frac{1}{4\sqrt{3}}(dx^{24} + 2dx^{35}), \\ \beta_{13}|_{\mathcal{S}^6} &= \frac{1}{4}dx^{25}, & \beta_{14}|_{\mathcal{S}^6} &= \frac{1}{4\sqrt{3}}(-dx^{25} + 2dx^{34}). \end{aligned}$$

Using these explicit formulae, we consider the evolution equation (2.11) subject to the initial constraint (2.12) for the ansatz (1.1). If we look at the e_5 component of A' , $*_r(dA \wedge r^4\psi)$ and $*_r(\frac{1}{2}[A, A] \wedge r^4\psi)$ at x we see that they are

$$\frac{1}{4}\frac{d(ar)}{dr}dx^2, \quad \frac{1}{4}2adx^2, \quad \frac{1}{4}\frac{1}{6}a^2r dx^2,$$

respectively. This suggests that the evolution equation (2.11) is equivalent to the ODE:

$$\frac{da}{dr} = \frac{a}{r}(1 + \frac{1}{6}ar). \quad (2.13)$$

Performing similar computations for the remaining 13 components allow us to verify that this is indeed the case. It is also straightforward to check that the constraint (2.12) holds automatically at x , the G_2 -instanton condition for the ansatz (1.1) is equivalent to solving (2.13).

In conclusion, we have shown the following.

Proposition 2.5. *The ansatz (1.1) on \mathbb{R}^7 with $G = G_2$ yields a non-trivial G_2 -instanton provided*

$$a(r) = -12\frac{r}{r^2 + C},$$

where C is a positive constant. This is the basic G_2 -instanton on \mathbb{R}^7 . \square

The instantons on $G_2 \times \mathbb{R}^7$ arising from Proposition 2.5 are well-known to physicists, arising in the context of heterotic string theory (see, for example, [11, 13]). In these papers, a more general ansatz is used, allowing for other gauge groups, but the only explicit solution provided is the one given above. More recently [7], it has been shown that this basic instanton is locally unique, but it is not known whether it is globally unique.

Remark 2.6. From the equations (2.10), one sees that the G_2 -instanton condition is conformally invariant, in the sense that the space of G_2 -instantons are the same for all G_2 -structures defined by 3-forms which differ by multiplication by a positive function. Therefore, it is natural to ask whether one can use stereographic projection to take the basic G_2 -instanton and define a G_2 -instanton on S^7 (perhaps away from the North pole) for its standard (or some other) G_2 -structure. However, one finds that the G_2 -structure one obtains away from the North pole on S^7 is not a multiple of the standard G_2 -structure on S^7 , and in fact cannot be extended to the North pole as a G_2 -structure. Hence, the basic instanton on \mathbb{R}^7 cannot be lifted in any obvious way to S^7 , unlike what occurs for the analogous setting of \mathbb{R}^4 and S^4 .

2.3. On \mathbb{R}^8 . Finally, we come to the basic instanton on \mathbb{R}^8 . We may regard $\text{Spin}(7)$ as the $\text{GL}(8, \mathbb{R})$ -stabiliser of the closed 4-form

$$\Phi = \varphi \wedge dx^8 + *_7\varphi,$$

where φ is given by (2.7) and $*_7\varphi$ by (2.8).

Again, we wish to write this as a conical structure on the open set $\mathbb{R}^8 \setminus \{0\} \cong \mathbb{R}_+ \times S^7$:

$$\Phi = r^3 \varphi \wedge dr + r^4 *_7\varphi \quad (2.14)$$

where φ is a G_2 -structure on S^7 satisfying $d\varphi = -4*_7\varphi$, in accordance with the torsion-free condition.

It is well-known that the condition for a 2-form β to lie in $\mathfrak{spin}(7) \subset \Lambda^2(\mathbb{R}^8)^*$ is given by:

$$*(\Phi \wedge \beta) = -\beta. \quad (2.15)$$

Again, the link to the idea of anti-self-duality is clear from this condition.

For our connection ansatz (1.1) we could again substitute $\beta = \Omega$ in (2.15) to obtain the $\text{Spin}(7)$ -instanton condition. However, we may also rephrase the $\text{Spin}(7)$ -instanton condition for the family $A(r)$ of connections over S^7 in (1.1) as

$$A' = *_r(F_A \wedge r^4 *_7\varphi) \quad (2.16)$$

for details see [15]. Here we used the decomposition (2.14) and $*_r$ is the Hodge star on the 7-sphere of radius r .

In order to obtain a basis of $\mathfrak{spin}(7) \subset \mathfrak{so}(8)$, we proceed as follows. Using the elementary matrices $E_{k\ell}$, a choice of orthogonal basis of $\mathfrak{spin}(7) \subset \Lambda^2(\mathbb{R}^8)^*$ is:

$$\begin{aligned} e_1 &= -E_{12} + E_{21} - E_{47} + E_{74}, & e_2 &= \frac{1}{\sqrt{3}}(-E_{12} + E_{21} + E_{47} - E_{74} - 2E_{56} + 2E_{65}), \\ e_3 &= \frac{1}{\sqrt{6}}(E_{12} - E_{21} - 3E_{38} + 3E_{83} - E_{47} + E_{74} - E_{56} + E_{65}), \\ e_4 &= -E_{13} + E_{31} - E_{46} + E_{64}, & e_5 &= \frac{1}{\sqrt{3}}(E_{13} - E_{31} - E_{46} + E_{64} - 2E_{57} + 2E_{75}), \\ e_6 &= \frac{1}{\sqrt{6}}(-E_{13} + E_{31} - 3E_{28} + 3E_{82} + E_{46} - E_{64} - E_{57} + E_{75}), \\ e_7 &= E_{14} - E_{41} - E_{27} + E_{72}, & e_8 &= \frac{1}{\sqrt{3}}(E_{14} - E_{41} + E_{27} - E_{72} - 2E_{58} + 2E_{85}), \\ e_9 &= \frac{1}{\sqrt{6}}(E_{14} - E_{41} + E_{27} - E_{72} - 3E_{36} + 3E_{63} + E_{58} - E_{85}), \end{aligned}$$

$$\begin{aligned}
e_{10} &= -E_{15} + E_{51} - E_{48} + E_{84}, & e_{11} &= \frac{1}{\sqrt{3}}(E_{15} - E_{51} - 2E_{26} + 2E_{62} - E_{48} + E_{84}), \\
e_{12} &= \frac{1}{\sqrt{6}}(-E_{15} + E_{51} - E_{26} + E_{62} - 3E_{37} + 3E_{73} + E_{48} - E_{84}), \\
e_{13} &= E_{16} - E_{61} - E_{78} + E_{87}, & e_{14} &= \frac{1}{\sqrt{3}}(-E_{16} + E_{61} - 2E_{25} + 2E_{52} - E_{78} + E_{88}), \\
e_{15} &= \frac{1}{\sqrt{6}}(-E_{16} + E_{61} + E_{25} - E_{52} - 3E_{34} + 3E_{43} - E_{78} + E_{87}), \\
e_{16} &= -E_{17} + E_{71} - E_{68} + E_{86}, & e_{17} &= \frac{1}{\sqrt{3}}(-E_{17} + E_{71} - 2E_{24} + 2E_{42} + E_{68} - E_{86}), \\
e_{18} &= \frac{1}{\sqrt{6}}(E_{17} - E_{71} - E_{24} + E_{42} - 3E_{35} + 3E_{53} - E_{68} + E_{86}), \\
e_{19} &= E_{18} - E_{81} - E_{23} + E_{32}, & e_{20} &= \frac{1}{\sqrt{3}}(E_{18} - E_{81} - 2E_{45} + 2E_{54} + E_{23} - E_{32}), \\
e_{21} &= \frac{1}{\sqrt{6}}(E_{18} - E_{81} + E_{23} - E_{32} + E_{45} - E_{54} - 3E_{67} + 3E_{76}).
\end{aligned}$$

The corresponding basis of 2-forms is:

$$\begin{aligned}
\beta_1 &= dx^{12} + dx^{47}, & \beta_2 &= \frac{1}{\sqrt{3}}(dx^{12} - dx^{47} + 2dx^{56}), \\
\beta_3 &= \frac{1}{\sqrt{6}}(-dx^{12} + 3dx^{38} + dx^{47} + dx^{56}), \\
\beta_4 &= dx^{13} + dx^{46}, & \beta_5 &= \frac{1}{\sqrt{3}}(-dx^{13} + dx^{46} + 2dx^{57}), \\
\beta_6 &= \frac{1}{\sqrt{6}}(dx^{13} + 3dx^{28} - dx^{46} + dx^{57}), \\
\beta_7 &= -dx^{14} + dx^{27}, & \beta_8 &= -\frac{1}{\sqrt{3}}(dx^{14} + dx^{27} - 2dx^{58}), \\
\beta_9 &= \frac{1}{\sqrt{6}}(-dx^{14} - dx^{27} + 3dx^{36} - dx^{58}), \\
\beta_{10} &= dx^{15} + dx^{48}, & \beta_{11} &= \frac{1}{\sqrt{3}}(-dx^{15} + 2dx^{26} + dx^{48}), \\
\beta_{12} &= \frac{1}{\sqrt{6}}(dx^{15} + dx^{26} + 3dx^{37} - dx^{48}), \\
\beta_{13} &= -dx^{16} + dx^{78}, & \beta_{14} &= \frac{1}{\sqrt{3}}(dx^{16} + 2dx^{25} + dx^{78}), \\
\beta_{15} &= \frac{1}{\sqrt{6}}(dx^{16} - dx^{25} + 3dx^{34} + dx^{78}), \\
\beta_{16} &= dx^{17} + dx^{68}, & \beta_{17} &= \frac{1}{\sqrt{3}}(dx^{17} + 2dx^{24} - dx^{68}), \\
\beta_{18} &= \frac{1}{\sqrt{6}}(-dx^{17} + dx^{24} + 3dx^{35} + dx^{68}), \\
\beta_{19} &= -dx^{18} + dx^{23}, & \beta_{20} &= \frac{1}{\sqrt{3}}(-dx^{18} + 2dx^{45} - dx^{23}), \\
\beta_{21} &= -\frac{1}{\sqrt{6}}(dx^{18} + dx^{23} + dx^{45} - 3dx^{67}).
\end{aligned}$$

To illustrate computations, let us consider the point $x \in S^7$ with $x_8 = 1$ and $x_j = 0$ for all other j . We then have that, at x ,

$$\begin{aligned}
\alpha_1 &= 0 = \alpha_2, & \alpha_3 &= -\frac{\sqrt{3}}{2}dx^3, & \alpha_4 &= 0 = \alpha_5, \\
\alpha_6 &= -\frac{\sqrt{3}}{2}dx^2, & \alpha_7 &= 0, & \alpha_8 &= -\frac{2}{\sqrt{3}}dx^5, & \alpha_9 &= \frac{1}{\sqrt{6}}dx^5, \\
\alpha_{10} &= -dx^4, & \alpha_{11} &= -\frac{1}{\sqrt{3}}dx^4, & \alpha_{12} &= \frac{1}{\sqrt{6}}dx^4, \\
\alpha_{13} &= -dx^7, & \alpha_{14} &= -\frac{1}{\sqrt{3}}dx^7, & \alpha_{15} &= -\frac{1}{\sqrt{6}}dx^7, & \alpha_{16} &= -dx^6, \\
\alpha_{17} &= \frac{1}{\sqrt{3}}dx^6, & \alpha_{18} &= -\frac{1}{\sqrt{6}}dx^6, & \alpha_{19} &= dx^1, & \alpha_{20} &= \frac{1}{\sqrt{3}}dx^1, & \alpha_{21} &= \frac{1}{\sqrt{6}}dx^1.
\end{aligned}$$

In addition, we have

$$\begin{aligned}
\beta_1 &= dx^{12} + dx^{47}, & \beta_2 &= \frac{1}{\sqrt{3}}(dx^{12} - dx^{47} + 2dx^{56}), & \beta_3 &= \frac{1}{\sqrt{6}}(-dx^{12} + dx^{47} + dx^{56}), \\
\beta_4 &= dx^{13} + dx^{46}, & \beta_5 &= \frac{1}{\sqrt{3}}(-dx^{13} + dx^{46} + 2dx^{57}), & \beta_6 &= \frac{1}{\sqrt{6}}(dx^{13} - dx^{46} + dx^{57}),
\end{aligned}$$

$$\begin{aligned}
\beta_7 &= -dx^{14} + dx^{27}, & \beta_8 &= -\frac{1}{\sqrt{3}}(dx^{14} + dx^{27}), & \beta_9 &= \frac{1}{\sqrt{6}}(-dx^{14} - dx^{27} + 3dx^{36}), \\
\beta_{10} &= dx^{15}, & \beta_{11} &= \frac{1}{\sqrt{3}}(-dx^{15} + 2dx^{26}), & \beta_{12} &= \frac{1}{\sqrt{6}}(dx^{15} + dx^{26} + 3dx^{37}), \\
\beta_{13} &= -dx^{16}, & \beta_{14} &= \frac{1}{\sqrt{3}}(dx^{16} + 2dx^{25}), & \beta_{15} &= \frac{1}{\sqrt{6}}(dx^{16} - dx^{25} + 3dx^{34}), \\
\beta_{16} &= dx^{17}, & \beta_{17} &= \frac{1}{\sqrt{3}}(dx^{17} + 2dx^{24}), & \beta_{18} &= \frac{1}{\sqrt{6}}(-dx^{17} + dx^{24} + 3dx^{35}), \\
\beta_{19} &= dx^{23}, & \beta_{20} &= \frac{1}{\sqrt{3}}(2dx^{45} - dx^{23}), & \beta_{21} &= -\frac{1}{\sqrt{6}}(dx^{23} + dx^{45} - 3dx^{67}).
\end{aligned}$$

We now consider the evolution equation (2.16) for our connection ansatz (1.1). Focusing on the coefficient of e_3 , we have that the projection of A' , $*_7(dA \wedge r^4 *_7 \varphi)$, $*_7(\frac{1}{2}[A, A] \wedge r^4 *_7 \varphi)$ are given by

$$-\frac{\sqrt{3}}{2} \frac{d(ra)}{dr} dx^3, \quad -2\frac{\sqrt{3}}{2} a(r) dx^3, \quad -3\frac{\sqrt{3}}{2} ra^2 dx^3,$$

respectively. These considerations suggest that the Spin(7)-instanton evolution equations (2.16) for the ansatz (1.1) is:

$$\frac{da}{dr} = \frac{a}{r}(1 + 3ar).$$

Indeed, it can be checked that the equations obtained from the remaining 20 components agree with this ODE.

Summing up, we have shown:

Proposition 2.7. *The ansatz (1.1) on \mathbb{R}^8 with $G = \text{Spin}(7)$ yields a non-trivial Spin(7)-instanton provided*

$$a(r) = -\frac{2}{3} \frac{r}{r^2 + C},$$

where C is a positive constant. This is the basic Spin(7)-instanton on \mathbb{R}^8 . \square

As for the G_2 -case, the corresponding family of instantons on $\text{Spin}(7) \times \mathbb{R}^8$ has been known to physicists for decades (see, for example, [8, 9]). Moreover, in this case a more general ansatz is used in the aforementioned papers, but only one explicit solution is provided, agreeing with the above. No uniqueness result is yet known for the basic instanton on \mathbb{R}^8 : even local uniqueness is currently open.

Remark 2.8. Just as for the G_2 case, the Spin(7)-instanton condition is conformally invariant, in that it does not change when multiplying the defining 4-form for the Spin(7)-structure by a positive function. However, here the possibility of lifting the basic instanton on \mathbb{R}^8 to \mathcal{S}^8 does not even arise as there is no Spin(7)-structure on \mathcal{S}^8 . This follows since the 8-sphere has non-vanishing Euler characteristic, χ , but vanishing Pontryagin classes, p_i , and so cannot satisfy the condition

$$8\chi = 4p_2 - p_1^2,$$

which is the necessary and sufficient condition for an orientable spin 8-manifold to admit a Spin(7)-structure (cf. [10]).

3. NAHM'S EQUATIONS AND QUATERNIONIC STRUCTURES

Prompted by Remark 1.3, we will now turn to an alternative ansatz that has appeared in the physics literature (see [13] and the references therein). This will allow us to construct families of explicit instantons on $\mathbb{H}^n = \mathbb{R}^{4n}$, essentially rephrasing [13, Sec. 5.1, Proposition 1] in index free notation. In contrast with the basic instanton constructions, the gauge group K is not imposed to be the same as the internal symmetry group G .

In this section, we let (J^1, J^2, J^3) be the standard triple of complex structures on $\mathbb{H}^n = \mathbb{R}^{4n}$ and let $\{k_j\}$ be any basis of the Lie algebra \mathfrak{k} of K .

3.1. Nahm's equations and instantons on \mathbb{H} . Consider first the task of finding an instanton on $\mathcal{U} \subset \mathbb{R}^4$ with given gauge group K . For a scalar function f defined on \mathcal{U} , we make the following connection ansatz

$$\omega = \sum_{j=1}^3 T_j(f) J^j(df), \quad (3.1)$$

where

$$T_i(s) = \sum_{j=1}^{\dim \mathfrak{k}} T_i^j(s) k_j, \quad i = 1, 2, 3,$$

is a \mathfrak{k} -valued function defined on an interval $I \subset \mathbb{R}$, containing the range of f . For this ansatz, we have the following recipe for constructing self-dual and instantons, slightly improving the statement of [13, Sec 6.1, Prop. 2]:

Proposition 3.1. *Consider a connection ω on a principal K -bundle over \mathbb{R}^4 given by the ansatz (3.1). If $T = (T_1, T_2, T_3)$, $T_i: I \rightarrow \mathfrak{k}$, satisfies Nahm's equations,*

$$T_i'(s) = [T_j(s), T_k(s)], \quad (3.2)$$

for each cyclic permutation $(ijk) = (123)$. Then:

(i) ω is self-dual if and only if f satisfies the system:

$$\frac{\partial^2 f}{\partial x_{pp}^2} = \frac{\partial^2 f}{\partial x_{qq}^2} \quad \text{and} \quad \frac{\partial^2 f}{\partial x_p \partial x_q} = 0, \quad (3.3)$$

for all $1 \leq p < q \leq 4$.

(ii) $-\omega$ is anti-self-dual if and only if f is harmonic.

Proof. The curvature of ω is given by

$$\Omega = \sum_{i=1}^3 \left(T_i'(f) df \wedge J^i(df) + T_i(f) d(J^i(df)) + [T_j(f), T_k(f)] J^j(df) \wedge J^k(df) \right), \quad (3.4)$$

where, as above, $(ijk) = (123)$ denotes a cyclic permutation of (123) . To proceed, we note that more explicitly,

$$\begin{aligned} J^1(df) &= -f_2 dx_1 + f_1 dx_2 - f_4 dx_3 + f_3 dx_4, \\ J^2(df) &= -f_3 dx_1 + f_4 dx_2 + f_1 dx_3 - f_2 dx_4, \\ J^3(df) &= -f_4 dx_1 - f_3 dx_2 + f_2 dx_3 + f_1 dx_4. \end{aligned}$$

Here we have used the notation $df = \sum_{j=1}^4 f_j dx_j$, where $f_j = \partial f / \partial x_j$. In particular, we have that

$$*df \wedge J^i(df) = J^j(df) \wedge J^k(df),$$

with $(ijk) = (123)$. Therefore, if T satisfies Nahm's equations, then self-duality becomes equivalent to the condition

$$d(J^i(df)) = *d(J^i(df)), \quad i = 1, 2, 3,$$

which, more explicitly, can be phrased in terms of the system (3.3). Similarly, anti-selfduality for $-\omega$ reads,

$$d(J^i(df)) = -*d(J^i(df)), \quad i = 1, 2, 3,$$

which is equivalent to requiring that f is harmonic. \square

Remark 3.2. A well-known solution, $T_i: (0, \infty) \rightarrow \mathfrak{k}$, to Nahm's equations, sometimes referred to as the Nahm pole solution, takes the form:

$$T_i(s) = \frac{1}{2s} k_i, \quad i = 1, 2, 3,$$

where k_1, k_2, k_3 span an $\mathfrak{su}(2)$ -subalgebra of \mathfrak{k} and satisfy the commutation relations $[k_1, k_2] = -2k_3$, etc. Recently, this solution has played a role in the study of the Kapustin-Witten equations (see, e.g., [18, 12]).

Proposition 3.1 allows us to produce families of explicit instantons on \mathbb{R}^4 . Some basic examples are given below.

Example 3.3. Given a solution of Nahm's equations and real constants A, B_j, C , any function of the form

$$f(\mathbf{x}) = A\mathbf{x} \cdot \mathbf{x} + \sum_{j=1}^4 B_j x_j + C$$

produces an SD instanton on \mathbb{R}^4 .

Example 3.4. For the ASD case, we can produce instantons on \mathbb{R}^4 starting from a solution to Nahm's equation and any function of the form

$$\begin{aligned} f(\mathbf{x}) = & A(\mathbf{x} \cdot \mathbf{x})^{-1} + A_1 x_1^2 + A_2 x_2^2 + A_3 x_3^2 - (A_1 + A_2 + A_3) x_4^2 \\ & + \sum_{1 \leq i < j \leq 4} B_{ij} x_i x_j + \sum_{j=1}^4 B_j x_j + C, \end{aligned}$$

with A, A_i, B_{ij}, B_i and C fixed real numbers; obviously, f is only defined away from zero if $A \neq 0$.

3.2. Quaternionic instantons on \mathbb{H}^n . The connection ansatz (3.1) can be used, more generally, for $\mathbb{H}^n = \mathbb{R}^{4n}$. The associated curvature 2-form is still given by (3.4), where, generalising the notation used above, we have:

$$\begin{aligned} J^1(df) &= \sum_{j=0}^{n-1} -f_{4j+2} dx_{4j+1} + f_{4j+1} dx_{4j+2} - f_{4j+4} dx_{4j+3} + f_{4j+3} dx_{4j+4}, \\ J^2(df) &= \sum_{j=0}^{n-1} -f_{4j+3} dx_{4j+1} + f_{4j+4} dx_{4j+2} + f_{4j+1} dx_{4j+3} - f_{4j+2} dx_{4j+4}, \\ J^3(df) &= \sum_{j=0}^{n-1} -f_{4j+4} dx_{4j+1} - f_{4j+3} dx_{4j+2} + f_{4j+2} dx_{4j+3} + f_{4j+1} dx_{4j+4}. \end{aligned}$$

In order to generalise the notion of ASD instantons (see, for example, [19]), we introduce the "fundamental" 4-form, Σ , on \mathbb{R}^{4n} given by

$$\Sigma = \frac{1}{2}(\sigma_1^2 + \sigma_2^2 + \sigma_3^2),$$

where

$$\begin{aligned} \sigma_1 &= \sum_{j=0}^{n-1} dx_{4j+1,4j+2} + dx_{4j+3,4j+4}, & \sigma_2 &= \sum_{j=0}^{n-1} dx_{4j+1,4j+3} + dx_{4j+4,4j+2}, \\ \sigma_3 &= \sum_{j=0}^{n-1} dx_{4j+1,4j+4} + dx_{4j+2,4j+3}. \end{aligned}$$

In these terms, we have the following definition.

Definition 3.5. A generalisation of ASD instantons on \mathbb{H}^n are given by connections ω whose curvature Ω satisfies:

$$\Omega \wedge \Sigma^{n-1} = c_1 * \Omega, \quad (3.5)$$

where $c_1 = -(2n-1)!/2^{n-1}$. In this case, the curvature 2-form takes values in $\mathfrak{sp}(1) \oplus \mathfrak{sp}(n) \subset \mathfrak{sp}(n+1)$.

Proposition 3.1 generalises to higher dimensional flat quaternionic space, \mathbb{H}^n , as follows.

Proposition 3.6. Consider a connection on a principal K -bundle over \mathbb{H}^n given by the negative, $-\omega$, of the ansatz (3.1). If $T = (T_1, T_2, T_3)$, $T_i: I \rightarrow \mathfrak{k}$, satisfies Nahm's equations (3.2), then $-\omega$ is anti-self-dual if and only if f satisfies the second order equations:

$$d(J^i(df)) \wedge \Sigma^{n-1} = c_1 * d(J^i(df)), \quad i = 1, 2, 3, \quad (3.6)$$

where $c_1 = -(2n-1)!/2^{n-1}$. \square

As in the 4-dimensional case, it is easy to construct families of explicit solutions as illustrated below for \mathbb{H}^2 .

Example 3.7. In higher dimensions, the conditions (3.6) are more involved than for \mathbb{R}^4 . For example, in the case of \mathbb{H}^2 this system of second order PDEs for f can be phrased as:

$$\begin{aligned} \frac{\partial^2 f}{\partial x_1^2} + \frac{\partial^2 f}{\partial x_2^2} + \frac{\partial^2 f}{\partial x_3^2} + \frac{\partial^2 f}{\partial x_4^2} &= 0 = \frac{\partial^2 f}{\partial x_5^2} + \frac{\partial^2 f}{\partial x_6^2} + \frac{\partial^2 f}{\partial x_7^2} + \frac{\partial^2 f}{\partial x_8^2} \\ \frac{\partial^2 f}{\partial x_1 \partial x_5} + \frac{\partial^2 f}{\partial x_2 \partial x_6} + \frac{\partial^2 f}{\partial x_3 \partial x_7} + \frac{\partial^2 f}{\partial x_4 \partial x_8} &= 0 = \frac{\partial^2 f}{\partial x_1 \partial x_6} - \frac{\partial^2 f}{\partial x_2 \partial x_5} - \frac{\partial^2 f}{\partial x_3 \partial x_8} + \frac{\partial^2 f}{\partial x_4 \partial x_7} \\ \frac{\partial^2 f}{\partial x_1 \partial x_7} + \frac{\partial^2 f}{\partial x_2 \partial x_8} - \frac{\partial^2 f}{\partial x_3 \partial x_5} - \frac{\partial^2 f}{\partial x_4 \partial x_6} &= 0 = \frac{\partial^2 f}{\partial x_1 \partial x_8} - \frac{\partial^2 f}{\partial x_2 \partial x_7} + \frac{\partial^2 f}{\partial x_3 \partial x_6} - \frac{\partial^2 f}{\partial x_4 \partial x_5}. \end{aligned}$$

From these equations, it follows that, with any solution of Nahm's equation, any function of the form

$$\begin{aligned} f(\mathbf{x}) &= f_1(\mathbf{x}) + f_2(\mathbf{x}) + C_1 x_1 x_5 + C_2 x_2 x_6 + C_3 x_3 x_7 - (C_1 + C_2 + C_3) x_4 x_8 \\ &\quad + D_1 x_1 x_6 + D_2 x_2 x_5 + D_3 x_3 x_8 - (D_1 - D_2 - D_3) x_4 x_7 \\ &\quad + E_1 x_1 x_7 + E_2 x_2 x_8 + E_3 x_3 x_5 + (E_1 + E_2 - E_3) x_4 x_6 \\ &\quad + F_1 x_1 x_8 + F_2 x_2 x_7 + F_3 x_3 x_6 + (F_1 - F_2 + F_3) x_4 x_5 \end{aligned}$$

where f_1 and f_2 are solutions of the type of Example 3.4 restricted to the J_i invariant subspaces $\{x_1, x_2, x_3, x_4\}$ and $\{x_5, x_6, x_7, x_8\}$, respectively, and C_i, D_i, E_i and F_i are constants.

Remark 3.8. In view of Proposition 3.6, it seems natural to ask for a generalisation of SD instantons on \mathbb{H}^n , using the ansatz (3.4). This would be a connection ω whose curvature Ω satisfies:

$$\Omega \wedge \Sigma^{n-1} = c_2 * \Omega, \quad (3.7)$$

where $c_2 = (2n+1)!/(6n2^{n-1})$. This corresponds to the curvature 2-form taking values in $\mathfrak{sp}(1) \subset \mathfrak{sp}(1) \oplus \mathfrak{sp}(n)$.

Direct computations show that the part of Proposition 3.1 concerning SD instantons does not generalise to higher dimensional \mathbb{H}^n .

In a sense, it is slightly misleading referring to the above instantons as being "quaternionic", since we are dealing with a flat hyperkähler structure on \mathbb{H}^n rather than a genuine quaternionic structure. Explicit examples of quaternionic ASD instantons are known, e.g., on $\mathbb{H}P^n$, cf. [3]; note that the authors of this paper use the opposite terminology for quaternionic ASD and SD instantons, respectively. These latter examples were revisited more recently in [4] in the context of hyperkähler cones.

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